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Continental interior and edge breakup at convergent margins induced by subduction direction reversal: A numerical modeling study applied to the South China Sea margin

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Abstract: The dynamics of continental breakup at convergent margins have been described as the results of back-arc opening caused by slab rollback or drag force induced by subduction direction reversal. Although the rollback hypothesis has been intensively studied, our understanding of the consequence of subduction direction reversal remains limited. Using thermo-mechanical modeling based on constraints from the South China Sea (SCS) region, we investigate how subduction direction reversal controls the breakup of convergent margins. The numerical results show that two distinct breakup modes, namely continental interior and edge breakup ("edge" refer to continent above the plate boundary interface), may develop depending on the "maturity" of the convergent margin and the age of the oceanic lithosphere. For a slab age of ~15-~45 Ma, increasing the duration of subduction promotes the continental interior breakup mode, where a large block of the continental material is separated from the overriding plate. In contrast, the continental edge breakup mode develops when the subduction is a short-duration event, and in this mode a wide zone of less continuous continental fragments and tearing of the subducted slab occur. These two modes are consistent with the interior (relic late Mesozoic arc) and edge (relic fore-arc) rifting characteristics in the western and eastern SCS margin, suggesting that variation in the northwest-directed subduction duration of the Proto-

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SCS might be a reason for the differential breakup locus along the strike of the SCS margin. Besides, a two-segment trench associated with the northwest-directed subduction is implied in the present-day SCS region.

Keywords: Continental breakup; Convergent margins; Edge breakup; Subduction direction reversal; Proto-South China Sea; Numerical modeling

1. Introduction

The South China Sea (SCS) is a typical extensional marginal sea basin similar to those found throughout the western Pacific. Despite structural similarities, the SCS also exhibits several distinctive features, e.g., along-strike variation in the breakup locus (e.g., *Briais et al., 1993; Huchon et al., 2001; Cullen et al., 2010; Li et al., 2012; Franke et al., 2014; Ding et al., 2016*). It is generally believed that before the opening of the SCS, a late Mesozoic ancient subduction zone existed along the SCS margin (e.g., *Li and Li., 2007; Xia and Zhao, 2014; Xu et al., 2016; Li et al., 2018a; Ye et al., 2018; Li et al., 2019a*). By recognizing the imprint of the late Mesozoic arc on the present-day SCS margin using petrological evidence and magnetic anomaly, *Li et al., 2018a* further inferred that the Cenozoic opening of the SCS basin broke the relic arc region (now continental interior) in the west and the relic forearc region (now continental edge) in the east (*Fig.1a*). In the west, the Dangerous Grounds block (Nansha block), which has an average crust thickness of ~24 km and width of ~100-300 km (*Fig. 1b*), was separated from South China and migrated southward to Borneo during 25-20.5 Ma (*Qiu et al., 2011; Barckhausen et al., 2014*). In contrast, the opening of the SCS in the east occurred

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at the thin continental edge and (*Fig. 1c*), as a result, the frontal part of the edge was split from the main continent and drifted southward since the Oligocene, finally forming the north Palawan (*Holloway, 1982; Maruyama et al., 1997; Zamoras and Matsuoka, 2004; Li et al., 2018a; Wang et al., 2019; Zhao et al., 2020; Zhang et al., 2020*). Petrological analyses of sedimentary rocks revealed a forearc-related origin of the north Palawan, which was filled by subduction-accretion complexes and juxtaposed chert-clastic sequences in an imbricate manner (*Isozaki et al., 1988; Marquez et al., 2006; Tumanda, 1991*). Another remarkable feature of the eastern SCS region is slab tearing. Many recent tomographic studies showed that an oceanic slab remnant (referred as the proto-SCS north slab) existed in the east, lying within the asthenosphere of the SCS and exhibiting a fast P-wave velocity anomaly (*Fig. 1d*) (e.g., *Hall and Spakman, 2015; Huang et al., 2015; Zahirovic et al., 2016; Wu and Suppe, 2018*). A similar but smaller slab remnant was also recognized by interpreting the wide-angle seismic profiles (*Zhou et al., 2006*). These slab remnants were possibly the results of a slab-tear event that occurred during the rifting of the continental edge (*Wu and Suppe, 2018; Sun et al., 2019*).

Previous studies have investigated the causes of continental rifting through analyzing geological records, geophysical data, and modeling results (e.g., *Turcotte et al., 1983; Buck, 1991; Bercovici, 2003; Péron-Pinvidic et al., 2009; Mohn et al., 2012; Coltice et al., 2019;*). Some of their end-members emphasized the role of plume-lithospheric interaction (e.g., *Behn et al., 2004; Phillips et al., 2005; Cande and Stegman, 2011; Husson, 2012; Yoshida and Hamano, 2015; Koptev et al., 2015, 2019; Jolivet et al., 2018; Mondy et al., 2018*), while others related the continental rifting to the tensional stresses generated at the plate boundary (e.g., *Huismans and Beaumont, 2003; Pérez-Gussinyé et al., 2006; Marotta et al., 2009; Choi et al.,*

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2013; Liao et al., 2013; Naliboff et al., 2017; Le Pourhiet et al., 2018, etc.). A series of numerical modeling studies further showed that, in both situations, tectonic loading, rollback process, pre-existing weak structure, extension rate, crustal rheology, subduction depth, and thermal structure can control the first-order patterns of continental breakup (e.g., *Pérez-Gussinyé et al., 2003; Lavier and Manatschal, 2006; Burov and Watts, 2006; Nemcok et al., 2013; Gueydan and Précigout, 2014; Svartman Dias et al., 2015; Dal Zilio et al., 2017; Ros et al., 2017; Tetreault and Buiter, 2018; Duclaux et al., 2019*). Among these factors, slab rollback is the common dynamics of continental rifting at convergent margins (e.g., *Stern, 2002; Leng and Gurnis, 2011; Heuret and Lallemand, 2005; Sdrolias and Muller, 2006*). For example, most of the marginal basins in the western Pacific, such as the Mariana Trough and Okhotsk Basin, have been considered as back-arc basins associated with the rollback process that took place along the trenches, where the Pacific and the Eurasian plates converge (*Karig, 1971*). During rollback subduction, the rifting features may be different by considering the slab-slab interactions (e.g., *Mishin et al., 2008; Faccenna et al., 2018; Holt et al., 2018*), trench migration rates (e.g., *Becker et al., 2015*), plate properties (e.g., *Capitanio et al., 2010; Rodríguez-González et al., 2012; Crameri and Tackley, 2015*), subduction erosion (e.g., *Von Huene and Scholl, 1991; Keppie et al., 2009*), and plate velocities (e.g., *Wolf and Huismans, 2019*). These former studies focused solely on the consequence of slab rollback and emphasized its role in breaking the convergent margins. Less attention, however, has been paid to other mechanisms, such as drag-force caused by the subduction direction reversal that occurred in the SCS region.

In the widely accepted subduction-collision model, the opening of the SCS was viewed as

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a consequence of drag force associated with southeastward subduction of a proto-oceanic plate (usually called Proto-SCS) underneath Borneo (e.g., *Hall, 2002; Hutchison, 2004*). The Proto-SCS (PSCS) was believed to be an oceanic crust which once occupied the area north of Borneo, wherein the present-day SCS is located (*Fig. 2a*). The tectonic reconstruction revealed that the subduction of the PSCS started along the South China margin and later shifted towards the southeast in the Eocene (~45 Ma) (*Hall, 2002; Wu and Suppe, 2018*) (*Fig. 2c and 2d*). The southeast-directed subduction invoked extension on the South China margin, which ultimately separated the Dangerous Grounds from the main continental body and created the SCS in between. The Dangerous Grounds was then dragged southward and finally collided with Borneo in the middle Miocene, which is regarded as the end of the PSCS subduction (e.g., *Holloway, 1982*) (*Fig. 1b* and *Fig. 2f*). The remnant slabs of the PSCS (north and south) are now trapped in the asthenosphere beneath the South China and Palawan, respectively (*Fid. 2f*) (*Fan et al., 2017; Wu and Suppe, 2018; Li et al., 2018a*). Although the geological processes have been discussed previously, there still lacks a quantitative test on this hypothesis. It remains unclear whether the subduction direction reversal can cause the continental breakup and, if so, how does this process break the continental edge or interior? Several numerical studies have been performed to investigate the formation of the SCS through modeling the lithospheric extension (e.g., *Brune et al., 2016; Le Pourhiet et al., 2018; Li et al., 2019b*), but their models focused mainly on rifting process without considering the role of subduction direction reversal. It is necessary to conduct a further study to improve our understanding of the rifting process in a convergent tectonic system.

Here we attempt to investigate the effects of subduction direction reversal on the

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development of continental breakup using a series of 2-D thermo-mechanical models. With these numerical models, we will also test the hypothesis that the "maturity" of convergent margins (in terms of subduction duration) and age of the oceanic lithosphere are crucial for explaining the contrasting breakup types along the strike of the SCS margin.

2. Numerical Model Description

The simulations in this study are performed using a 2-D thermo-mechanical code modified after *Gerya (2010)*, solving the Stokes and heat transfer equations. More details about the governing equations, material parameters, and benchmarks used in the numerical experiments are in the supplementary material and example of 'viscoelastic-plastic subduction' in the reference book.

We design a model domain of 1100 km×500 km in the x- and y- dimensions with a non-uniform rectangular grid and randomly distributed markers. The Eulerian grid in the vicinity of the subduction zone has a higher resolution of 1 km × 1 km, whereas the rest of the model increases their grid spaces away from the high-resolution domain by a constant gradient of 1.1. In total ~1.6 million random Lagrangian markers are evenly distributed in the Eulerian cells.

Our initial setup is a standard subduction model that consists of two parts: (1) an upper continental plate and (2) a subducting slab. The continent plate is composed of an upper crust and a lower crust with a total thickness of 35 km, underlain by lithospheric mantle up to a depth of 85 km that was estimated from a regional S-velocity model (*Fig. 3a* and *3d*) (*An and Shi, 2006*). Following earlier similar studies (e.g., *Liao et al., 2017; Beaussier et al., 2019; Koptev et al., 2019*), the upper and lower continental crust in our numerical experiments are prescribed with quartzite and plagioclase rheology, respectively (*Shelton et al., 1981; Ranalli and Murphy,*

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1987). The mantle is represented by dry olivine rheology which is controlled by a combination of Power-law, diffusion and Peierls creep (*Chopra and Paterson, 1981; Kirby, 1983; Karato and Wu, 1993; Durham et al., 2009*) (*see Supplementary Material*). The oceanic crust is given by a 2-km-thick basalt-like oceanic upper crust and a 5-km-thick gabbro-like oceanic lower crust (e.g., *Turcotte and Schubert, 2002*) (*Fig. 3d*). We describe the brittle behavior of rocks by the classical Mohr-Coulomb failure criterion (see e.g., *Vermeer, 1990*). Strain softening is specified as a linear decrease of the friction angle with increasing accumulated strain (see, e.g., *Gerya, 2010*).

Statistical analyses of trenches in the western Pacific (*Lallemand et al., 2005*) and near-field seismic observations (*Zhu et al., 2019*) show an average slab dip of 28° at the shallow depth (*Fig. 3b*), so we initiate the subduction by prescribing a 28° dipping weak fracture zone between the continental and oceanic plates. This weak zone consists of mantle rocks with wet olivine rheology (*Chopra and Paterson, 1981; Ranalli, 1995*) and low plastic cohesion (1 MPa), thus facilitating the subduction along the channel (*Fig. 3d*). A 7 km-thick 'sticky air' ($\eta=10^{18}$ Pa·s, $\rho=1$ kg/m^3) layer is imposed on the model's upper surface to provide a free-surface-like condition that allows topographic evolution (*Crameri et al., 2012*). The details of the material parameters are illustrated in Table S1.

Our model has a constant temperature boundary condition at the top and bottom boundaries, while the left and right boundaries are thermally isolated. The initial thermal structure of the continental lithosphere increases vertically from 0 °C at the surface to 1380 °C at 85 km in a linear trend (*Fig. 3c*). The initial temperature of the asthenosphere mantle increases with depth at a gradient of 0.5 °C/km. The thermal structure of the oceanic plate is defined by a half-space

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cooling age (*Turcotte and Schubert, 2002*). However, the age of the PSCS is still in debate, and previous studies have argued that the PSCS was possibly a remnant of the Paleo-Pacific/Meso-Tethys, a Mesozoic oceanic embayment, a back-arc basin, or a marginal basin (e.g., *Taylor and Hayes, 1983; Holloway, 1982; Zhou et al., 2005*). Therefore, we vary the ages of the oceanic lithosphere (e.g., 10 Ma, 20 Ma, 30 Ma, 40 Ma, and 50 Ma) to explore their potential effects (*Fig. 3f*).

The mechanical boundary conditions are free slip at the top and left side boundaries. The bottom boundary is permeable in the vertical direction to satisfy an external/internal free slip boundary condition (*Burg and Gerya, 2005*). To represent the motion of the PSCS, we assign a constant convergent velocity followed by divergent velocity on the right side of the model domain. During the first 4 Myr (5 to 12 Myr also tested), we impose a constant convergence velocity of -3 cm/yr to initiate the leftward subduction (*Fig. 3d* and *3e*). To get a smooth transition from subduction to extension, we follow the approach of *Jourdon et al., 2019*, adopting a 2 Myr window between convergence and divergence to simulate a period of tectonic quiescence. That is, the velocity is zero from 4 Myr to 6 Myr. Then, after 6 Myr a constant divergence velocity of 1.5cm/yr is prescribed on the right side until the end of the simulation (*Fig. 3e*). More models are conducted by varying the subduction duration from 4 to 12 Myr to investigate the role of convergent margin "maturity", and hence that the extension starts at 6 to 14 Myr, respectively.

To differentiate the initial breakup locations, we divide the continental domain into two lithospheric units: the continental edge and interior, depending on whether the continent directly overlies the plate boundary interface (*Fig. 3d*). The continental edge is above the

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interface with a thin lithosphere, while the continental interior locates adjacent to the continental edge and is characterized by a stable and larger lithospheric thickness than that of the continental edge.

3. Results:

To test the hypotheses concerning the variable lithospheric age and subduction duration, several groups of numerical experiments are designed in this study. Each group includes a given lithospheric age and nine different subduction durations (4 to 12 Myr).

3.1 Group 1: lithospheric age of 10 Ma and variable subduction duration

Group 1 has a lithospheric age of 10 Ma. The results with subduction duration of 4, 6, 8, and 10 Myr are selected in *Fig. 4* (models **1-4**). All these models experience a steady-state subduction process during convergent margin evolution, and the slabs finally reach different depths at the end of the subduction stage.

During lithospheric extension, these models display a similar first-order deformation behavior. Moreover, at the end of the extension, the initial breakup occurs at the continental edge that specifies the leading edge of the upper continent overlying the plate boundary interface. For example, as revealed by model **1**, the divergent velocity exerts a negligible effect on the continent at 6 Myr, so that the continent does not involve in the extensional activity at the early stage of the lithospheric extension. At the same moment, the relative motion along the plate boundary interface accommodates almost the entire model's deformation (*Fig. 4a, at 6 Myr*). At 12 Myr, the oceanic plate returns to a shallow depth. The intense strain localization begins to occur at the continental edge and propagates into the underlying oceanic lithosphere (*Fig. 4a, at 12 Myr*). At 19 Myr, the continental edge eventually breaks into pieces along the

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shear zone. These newly formed continental fragments move together with the oceanic plate to the right side (*Fig. 4a, at 19 Myr*). Another distinctive characteristic is the tearing of the subducted slab. Therefore, a slab remnant can be observed below the overriding continental plate at the end of the modeling (*Fig. 4a, at 19 Myr*).

Also note that, in the long-period subduction cases (model **3** and **4**), the incipient breakup is delayed by ~9-15 Myr compared to that in the short-period subduction cases (model **1** and **2**). This addition time ultimately leads to the slab remnant sinking into the asthenosphere (*Fig. 4c, at 28 Myr* and *Fig. 4d, at 34 Myr*).

3.2 Group 2: lithospheric age of 20 Ma and variable subduction duration

Group 2 has a lithospheric age of 20 Ma. We also select the models with subduction durations of 4, 6, 8, and 10 Myr for comparison (*Fig. 5*, models **5-8**). The evolution of models in group 2 becomes quite different when the tectonic inversion starts. The major part of the strain localizes in the continental edge and the oceanic slab (models **5-7**), or in the continental interior (model **8**), depending on the subduction duration. As a result, two distinct continental breakup modes, namely continental edge and interior breakup, develop at the end of the experiment.

Similar to the results in group 1, models **5** -**7** exhibit the rift features of the continental edge breakup mode, e.g., relative motion between the plates, tearing of the subducted slab, and formation of the continental fragments. In contrast, model **8** shows another scenario in which no clear relative motion between the plates occurred during extension. Instead, the continent is subject to continuous stretching, and as a result, lithospheric-scale shear bands develop in the continental interior to accommodate most of the extensional strain (*Fig. 5d, at 20 Myr*). Such

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differences possibly stem from that the slab in model **8** has been subducted into greater depth, thus increasing the inter-plate contact area and exhibiting a strong plate coupling during the rifting stage. As extension progresses, the continent extends over a large scale and develops a typical necking in the lithospheric mantle, similar to those in previous studies (e.g., *Duretz et al., 2016; Chenin et al., 2017; Li et al., 2019*). At the end of the system's development, breakup initiates in the continental interior, which is approximately 200 km away from the trench, separating a microcontinent from the main continent body. The microcontinent migrates rightward together with the oceanic plate, while the rest of the continent remains in its original location forming a passive margin (*Fig. 5d, at 24 Myr*).

3.3 Group 3: lithospheric age of 30 Ma and variable subduction duration

Group 3 has an oceanic age of 30 Ma. Figure 6 summarizes the evolution of models with subduction durations of 7, 8, 9, and 10 Myr. The essential features of this group are similar to group 2. Group 3 also concludes that breakup can occur at the continental edge or within the interior, depending on the subduction duration.

It is also noteworthy that an increase in the lithospheric age can enhance the mechanical locking at the plate interface. Thus, compared to the earlier groups, group 3 requires an even shorter subduction duration to facilitate the development of interior breakup. For example, in contrast to group 2 where the model with subduction duration of 6 Myr is characterized by continental edge breakup (model **6**), here the 6 Myr induces the breakup of the continental interior (model **10**, *Fig. 6b*). This implies that, in the case of the oceanic age of 30 Ma, the continental edge breakup only occurs when the subduction ends earlier than 6 Myr (*Fig. 6a*).

4. Discussion

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4.1 Summary of numerical results

Our experiment results demonstrate that variations in the subduction duration when the reversal of subduction direction happened, may result in two contrasting continental breakup modes at the end of the rifting stage. In general, short-period subduction duration shows a weak coupling condition and, as a result, the relative motion along the plate boundary interface accommodates the model's deformation at the early extension stage. When the oceanic slab returns to a shallow depth, the deformation begins to localize at the continental edge and eventually breaks the continental edge (*Fig. 7b*). In the edge breakup mode, the extension also tears the subducted oceanic slab into two plates. The main oceanic plate moves rightward, whereas the remnant is trapped beneath the continental lithosphere (*Fig. 7b*). In contrast, the long-period subduction duration model exhibits a strong coupling between the subducting and overriding plates when the subduction direction reversal occurs. In these strongly coupled cases, deformation mainly localizes within the continent interior, which finally breaks up at the end of rifting (*Fig. 7c*).

As shown in *Fig. 7d*, we summarize the models in terms of breakup types as a function of subduction duration and age of the oceanic lithosphere. The results show that, for a given slab age of ~15-~45 Ma, the less the subduction duration is, the greater the likelihood of continental edge breakup will occur (shadow area in *Fig. 7d*). For example, subduction duration should be less than 8 Myr to induce edge breakup for a slab age of 20 Ma, otherwise the breakup will occur within the continental interior. Furthermore, as the slab age increases, the subduction duration required to induce the continental edge breakup is reduced. For example, the age of 20 Ma can promote the breakup of continental edge on the condition that the subduction ends

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earlier than 8 Myr, whereas the age of 30 Ma requires a subduction duration less than 4 Myr (shadow area in *Fig. 7d*).

The results also suggest that, once the slab age exceeds ~45 Ma, the breakup of the continental edge has never been observed because the mechanical plate coupling is sufficiently strong to resist extension. Yet the continental edge breakup always occurs when the age is younger than ~15 Ma due to weak plate coupling (*Fig. 7d*).

4.2 Models consistent with observations in the SCS continental margin

In the SCS continental margin, the relic Mesozoic volcanic arc was broken and extended over wide area in the west, whereas the arc in the east is less extended and kept its original configuration (*Li et al., 2018a*). These observations are consistent with our numerical results. In the interior breakup mode, strain mainly localized further away from the trench and along the continental interior during extension. As extension proceeded, the arc may be stretched over wide area and finally was split by the opening of oceanic basin (*Fig. 7c*). The western SCS margin followed the interior breakup mode and thus developed a wide range of arc in both conjugate margins of the southwestern sub-basin. By contrast, the eastern SCS margin was ruptured along the relic forearc region and, as a result, the late Mesozoic arc was rarely extended and was left in the Chinese side with its original configuration during the opening of SCS basin. Another important characteristic in the east is that a portion of the PSCS slab lay under the present-day SCS and exhibiting a high-velocity feature (*Fig. 1d*) (*Wu and Suppe, 2018*). These aforementioned characteristics are in agreement with our continental edge mode, where the breakup occurred along the forearc region and was accompanied by the tearing of the subducted slab (*Fig. 7b*).

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By reconstructing the plate boundary between Sundaland and the Dangerous Grounds, *Clift et al. (2008)* proposed that the PSCS was about one-thousand-kilometer-wide in the north-south direction (*Fig. 8*). Many previous studies believed that the PSCS has been subducted northwestward beneath the South China block and southeastward beneath the northern Borneo (*Hall, 2002; Wu and Suppe, 2018*) (*Fig. 2*). The southeast dipping subduction goes in a scissor-like fashion, in which a larger portion of the oceanic plate has been subducted beneath Sarawak than beneath Sabah (*Madon et al., 2013*) (*Fig. 8*). Therefore, the amount of the PSCS subducted beneath the South China block can be roughly estimated by subtracting the southeast subducted slab from the total width. The results shown in *Fig. 8* indicate that a larger portion of the PSCS is subducted beneath the western South China margin than that beneath the eastern South China margin. In other words, the western South China margin experiences a longer period of subduction compared to the eastern margin.

The variation in the subduction duration has a good correspondence with the differential breakup characteristics along the strike of the South China margin. As shown in *Fig.1* and *Fig. 8*, the eastern South China margin is characterized by the breakup of the edge region and experiences a relatively short period of the PSCS subduction. By contrast, the western margin has been broken within the interior and undergoes a relatively long period of subduction. These phenomena are consistent with our numerical results, where a long subduction duration induces the continental interior breakup, whereas a short duration causes the continental edge breakup. Therefore, our studies demonstrate that variation in the duration of northwest dipping subduction might be a cause for the contrasting breakup types along the strike of the SCS margin.

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4.3 Models implication for the northwest-directed subduction trench

Previous studies have revealed that the southeast-directed subduction trench locates at the southern SCS margin, but the northwest-directed subduction trench remains unclear. In combination with the numerical results and geophysical characteristics, this section will discuss the possible northwest-directed trench locations. As mentioned by previous studies (e.g., *Zhou et al., 2008*), the horizontal gradient of the Bouguer anomaly (HGBA) is sensitive to lateral material contrast and thus can illustrate the tectonic boundaries. In the HGBA map of the SCS region, a few high-amplitude anomaly belts are observed clearly parallel/subparallel to the margin (*Fig. 9a*). We number these belts and find that many of them are well correlated with the known plate boundaries in the SCS region. For example, belt 2 delineates the shape of the SCS oceanic basin, consistent with the continent-ocean boundaries. The Manila trench, where the SCS plate is subducting beneath the Philippine Sea plate, is delineated by the high-amplitude anomaly belt 3.

Belt 4 is located at the southern end of the Nansha Trough (NW Borneo trough) and corresponds to a relic trench, also known as the thrust front that is related to the southeast-directed subduction of the Proto-SCS beneath Borneo (*Hamilton, 1979; Franke et al., 2008*). The southeast-directed subduction ceased in the Early Miocene when the Dangerous Grounds collided with Borneo causing compressive deformation at the thrust front (e.g., *Hutchison, 1996; Hall, 2008*). As collision proceeded, the attenuated continental crust of the Dangerous Grounds went underneath Borneo. The southern Dangerous Grounds leading edge has been thrust beneath the Crocker-Palawan accretionary wedge, possibly as far to the onshore beneath Mount Kinabalu (MK in *Fig. 9a*) in Borneo (e.g. *Cottam et al., 2010; Cullen et al., 2010;*

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Franke et al., 2008; Hutchison and Vijayan, 2010; Steuer et al., 2013). Belt 5 well crosses Mount Kinabalu and overlaps the regions of the Crocker-Palawan complex, thus possibly denoting the southern boundary of the Dangerous Grounds (*Fig. 9b*). According to the evolution of our continental interior breakup mode, the Dangerous Grounds is separated from the main South China continent, and its southern boundary might be the site where the northwest-directed trench was ever located (*Fig. 9c*).

We also note that a distinct belt (named Belt 1) in the northeastern SCS margin does not correspond to any known plate boundaries, yet shows strong lateral material contrast, extending from the Central Uplift to the Yitongansha (*Fig. 9a*). Integrating numerical results with observations, we infer the belt probably due to the existence of slab remnant that is trapped beneath the continent during the continental edge breakup. The breakup of the continental edge tears the oceanic plate and causes the remnant staying in the subduction channel or sinking into the asthenosphere, thus causing a strong lateral material contrast (*Fig. 9c*). Our interpretation is also supported by other independent evidence. For instance, a remnant of the oceanic crust has been identified in the middle crust from seismic reflections in the northern SCS margin (e.g., *Huang et al., 2005; Zhou et al., 2006; Hu et al., 2008*). These identified crust remnants coincide with the gravity anomaly belt and thus provide additional support to our interpretation. In conclusion, the opening of the SCS basin has broken the continental interior in the west, moving the west part of the northwest-directed trench (Belt 5) southward to Borneo with the Dangerous Grounds. By contrast, the initial breakup site for the east SCS is the continental edge, and thus, the east part of the northwest-directed trench (Belt 1) remains in the northern SCS margin.

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4.4 Model limitations

Our thermo-mechanical experiments consider the southeastward subduction of the PSCS as the driving force that induces the continental breakup in the context of convergence. However, we only tested the effects of subduction duration and oceanic age. Other factors, such as surface process (e.g., *Olive et al., 2014; Erdős et al., 2014; Andrés-Martínez et al., 2019*), rheological heterogeneities (e.g., *Vogt et al., 2017*), plate motion direction (*Duclaux et al., 2019; Sun et al., 2020*), fluids (*Gerya and Meilick, 2011*), etc., might have also facilitated this breakup process, but investigating all their contributions are beyond the scope of this study. In the future, the effects of these factors on the evolution of the SCS should be investigated and evaluated.

Our simulations are here simplified using a 2-D process, which does not allow analyzing the 3-D effects. 3-D numerical simulations are necessary tools to investigate the lateral strain propagation during the model's evolution (e.g., *Menant et al., 2016; Beniest et al., 2017; Koptev et al., 2017*). However, as we focus on evolution from rift initiation to breakup and not on rift propagation, the use of a 2-D setup in our study is meaningful to investigate the conditions of creating a new rifted margin. In addition, the 2-D subduction model can generate large enough mechanical force to break the continent, in accordance with previous works (*Leng and Gurnis, 2011; Tan et al., 2012*). Although simplified, our present models are applicable to natural examples and thus can provide first-order estimates of continental breakup induced by subduction direction reversal.

5 Conclusions

In this study, we investigate how the subduction direction reversal controls the continental breakup styles using thermo-mechanical modeling based on the available geological and

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geophysical constraints from the SCS. The results show that the drag force induced by subduction direction reversal can invoke breakup of the South China continent, and two breakup modes, namely, continental interior and continental edge breakup, are developed depending on the slab age and subduction duration. The continental interior breakup mode is characterized by separating a large continental from the overriding plate. By contrast, the continental edge breakup mode develops a wide zone of less continuous continental fragments and tearing of the subducted slab. For a slab age of ~15-~45 Ma, the continental interior breakup mode requires a longer subduction duration than the edge breakup mode. The continental edge breakup always occurs when the slab age is younger than ~15 Ma, whereas the breakup only develops within the interior once the age exceeds ~45 Ma.

By comparing the modeling results to the final rifting characteristics of the SCS, our results suggest that the variation in the subduction duration might be a reason for the contrasting breakup locus along the strike of the SCS margin. That is, the eastern South China margin has experienced a relatively short-period PSCS subduction and thus is characterized by the breakup of the edge region, whereas the western South China margin undergoes a long period subduction and thus has been broken in the interior region.

A combination of HGBA and numerical modeling suggests that a two-segment northwest-directed trench exists in the present-day SCS region. The western segment has moved southward together with the Dangerous Grounds and now lies beneath the Crocker-Palawan accretionary wedge. The eastern segment is found in the northern SCS margin, extending from the Central Uplift to the Yitongansha.

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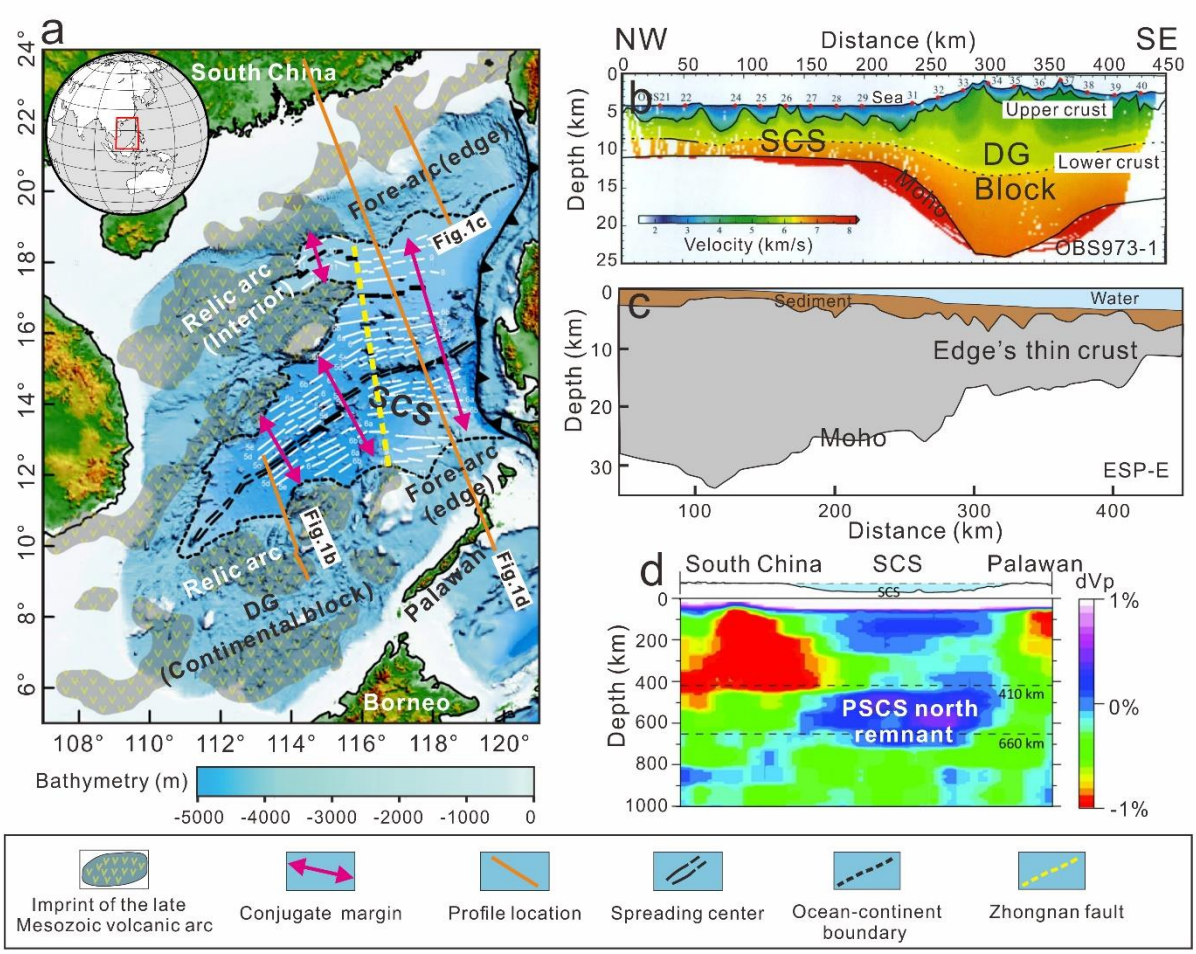

Figure 1. Observations in the South China Sea region. (a) Bathymetry map and the imprint of the late Mesozoic volcanic arc in the SCS region. The imprint of the late Mesozoic volcanic arc shows that the Cenozoic opening of the SCS basins has broken the continental interior in the west and the continental edge in the east (*Li et al., 2018a*). (b) and (c) Crustal velocity structures beneath the southwestern and northeastern margin (*Nissen et al., 1995; Qiu et al., 2011*). For more crustal structures in SCS, please refer to OBS2001 in *Zhao et al., 2010* and OBS2006 in *Wei et al., 2011*. (d) Tomographic section indicating a sub-horizontal fast slab anomaly (the PSCS north slab) under the east SCS (*Wu and Suppe, 2018; Sun et al., 2019*). DS: Dangerous Grounds (Nansha).

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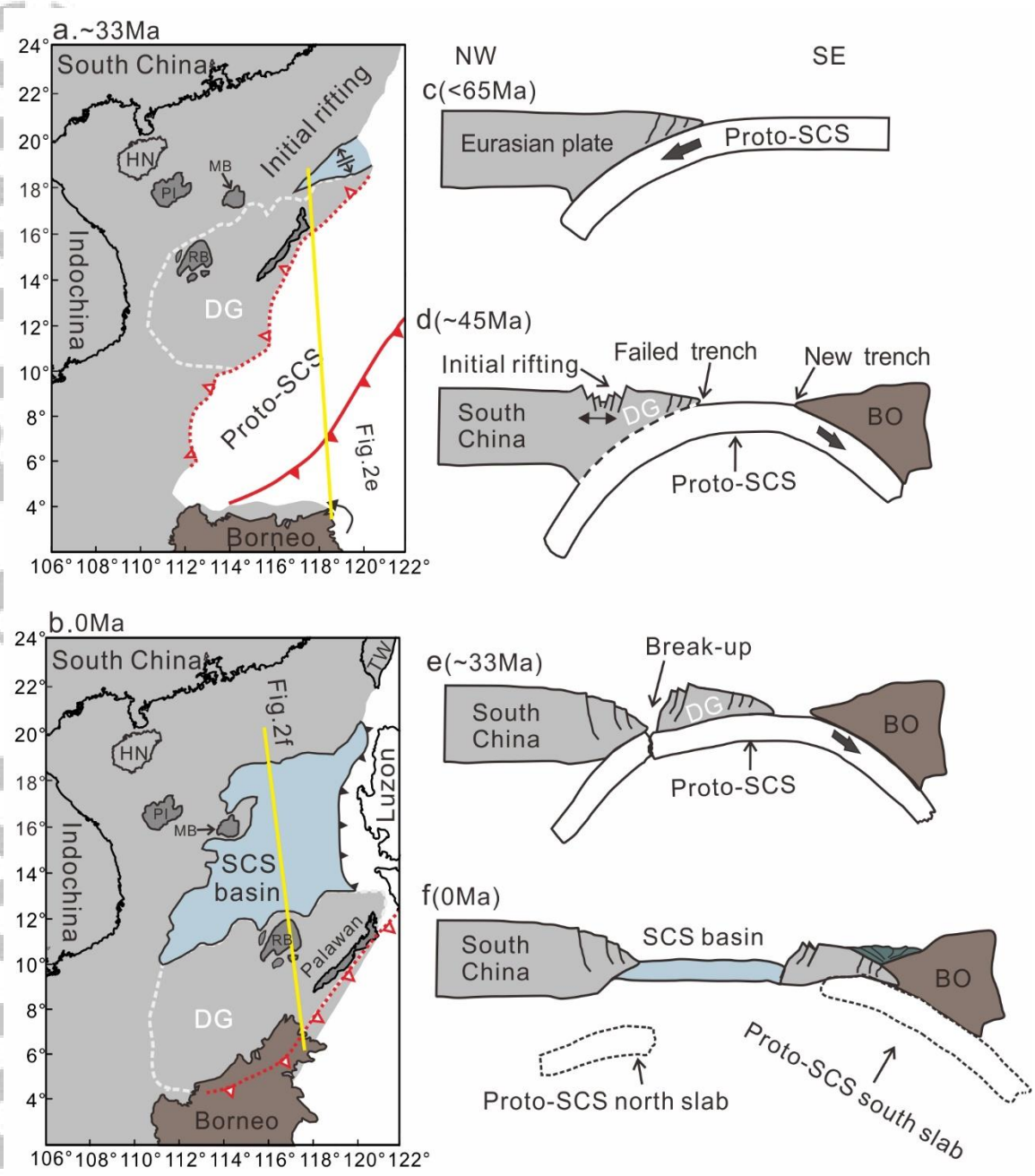

Figure 2. Plate reconstructions of the SCS region and the schematic of the process of subduction direction reversal (summarized after *Holloway, 1982; Pautot et al., 1986; Hall, 2002; Sun et al., 2009; Li et al., 2014; Wu and Suppe, 2018; Li et al., 2018b; Sun et al., 2019*). (HN: Hainan; MB: Zhongsha Bank (Macclesfield Bank); PI: Xisha Islands (Paracel Islands); RB: Reed Bank; BO: Borneo; DG: Nansha (Dangerous Grounds); PH: Philippine; PN: Palawan; SCS: South China Sea.)

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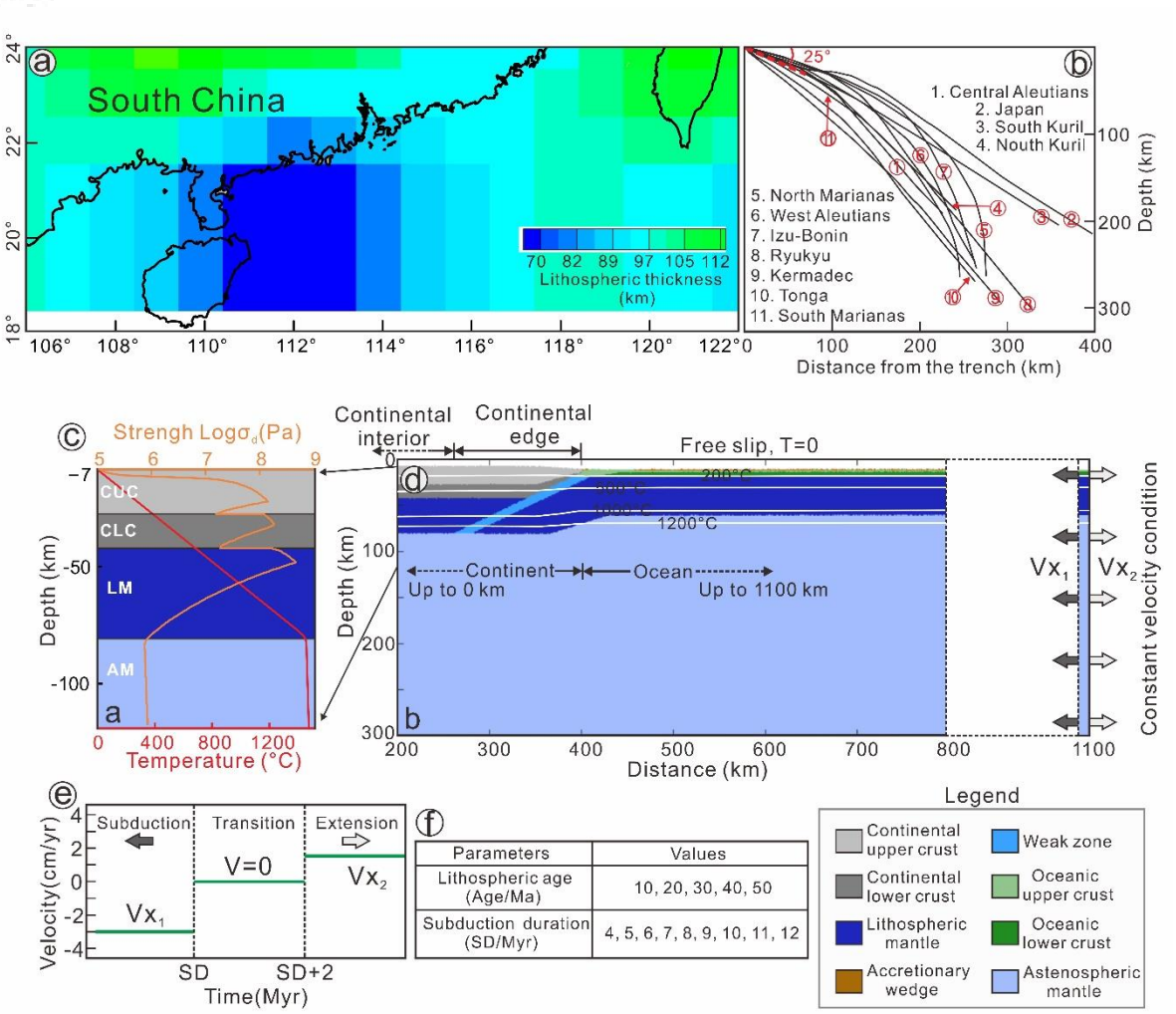

Figure 3. Initial model setup and constraints: (a) seismic-thermal lithosphere thickness of the South China block and surrounding SCS margin (*An and Shi, 2006*), (b) statistical analysis of subduction dips in the west Pacific (summarized after *Lallemand et al., 2005*), (c) yield strength profile (brown line) and geotherm (red line) of the continent, (d) enlarged 600 km × 300 km domain of the original 1100 km × 300 km model (continental edge and interior in the modes depending on whether the continent directly overlies the plate boundary interface), (e) variable velocity boundary conditions on the right side of the model domain, (f) key variable parameters of the numerical experiments. CUC: continental upper crust; CLC: continental lower crust; LM: lithospheric mantle; AM: asthenospheric mantle.

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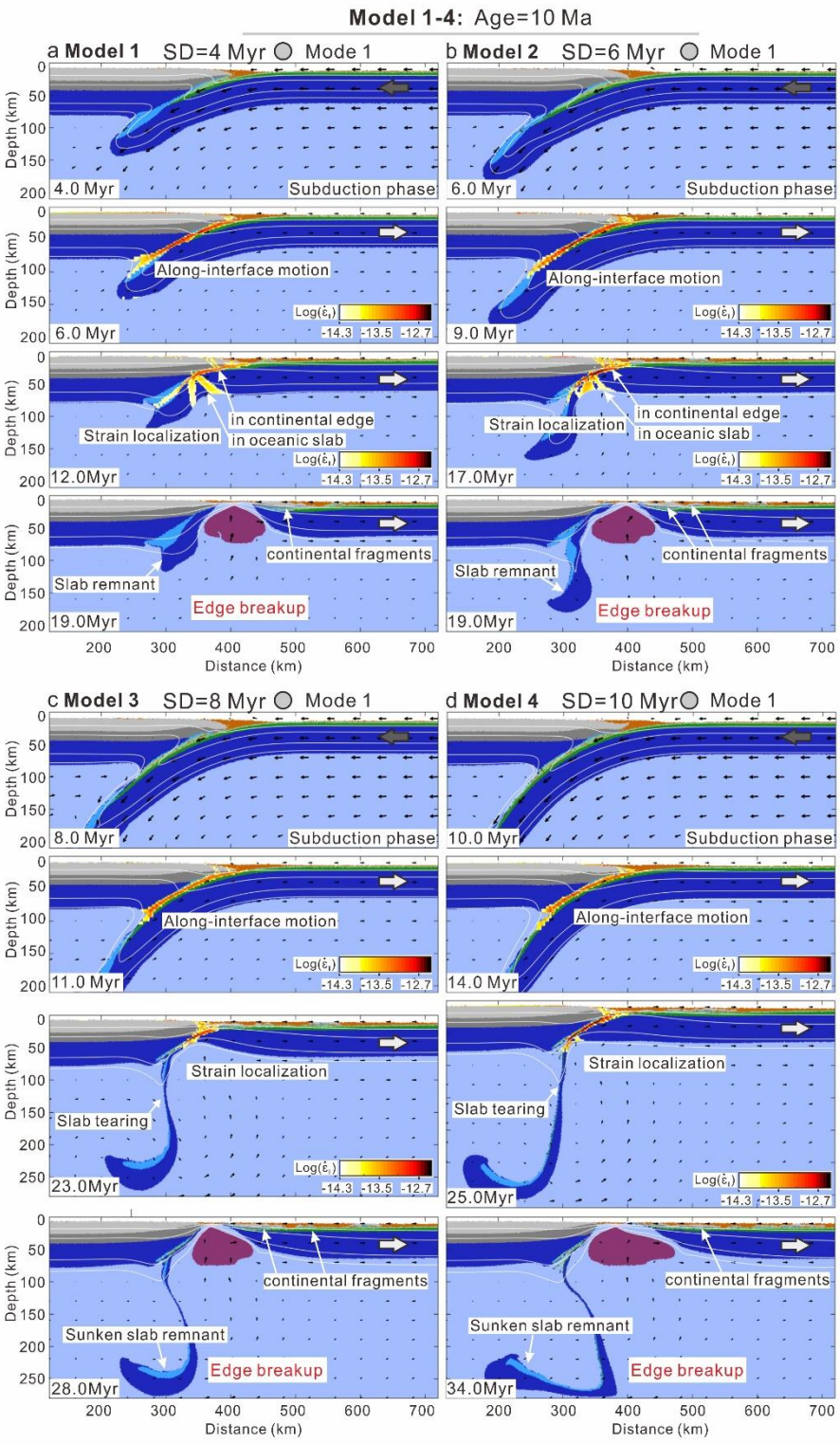

Figure 4. Evolution of group 1. Oceanic lithospheric age of 10 Ma and subduction duration (SD) of 4 (model 1), 6 (model 2), 8 (model 3), 10 Myr (model 4). The inset colors show the second invariant of the strain rate (as in the colorbar). White lines are isotherms in °C. Black arrows denote the calculated velocity field. Colors of rock types are as on Figure 2.

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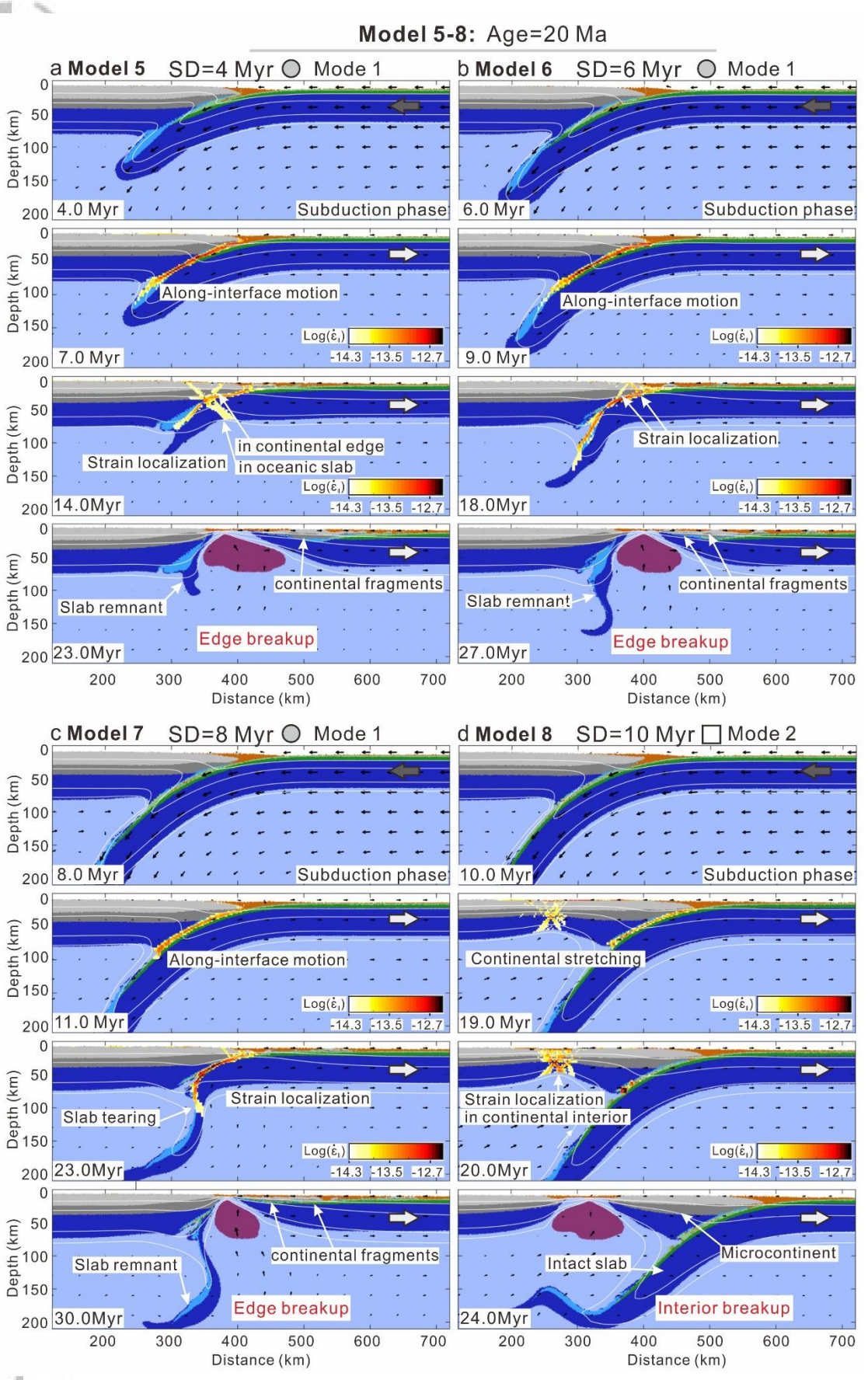

Figure 5. Evolution of group 2. Oceanic lithospheric age of 20 Ma and subduction duration

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(SD) of 4 (model 5), 6 (model 6), 8 (model 7), 10 Myr (model 8). The inset colors show the second invariant of the strain rate (as in the colorbar). White lines are isotherms in °C. Black arrows denote the calculated velocity field. Colors of rock types are as on Figure 2.

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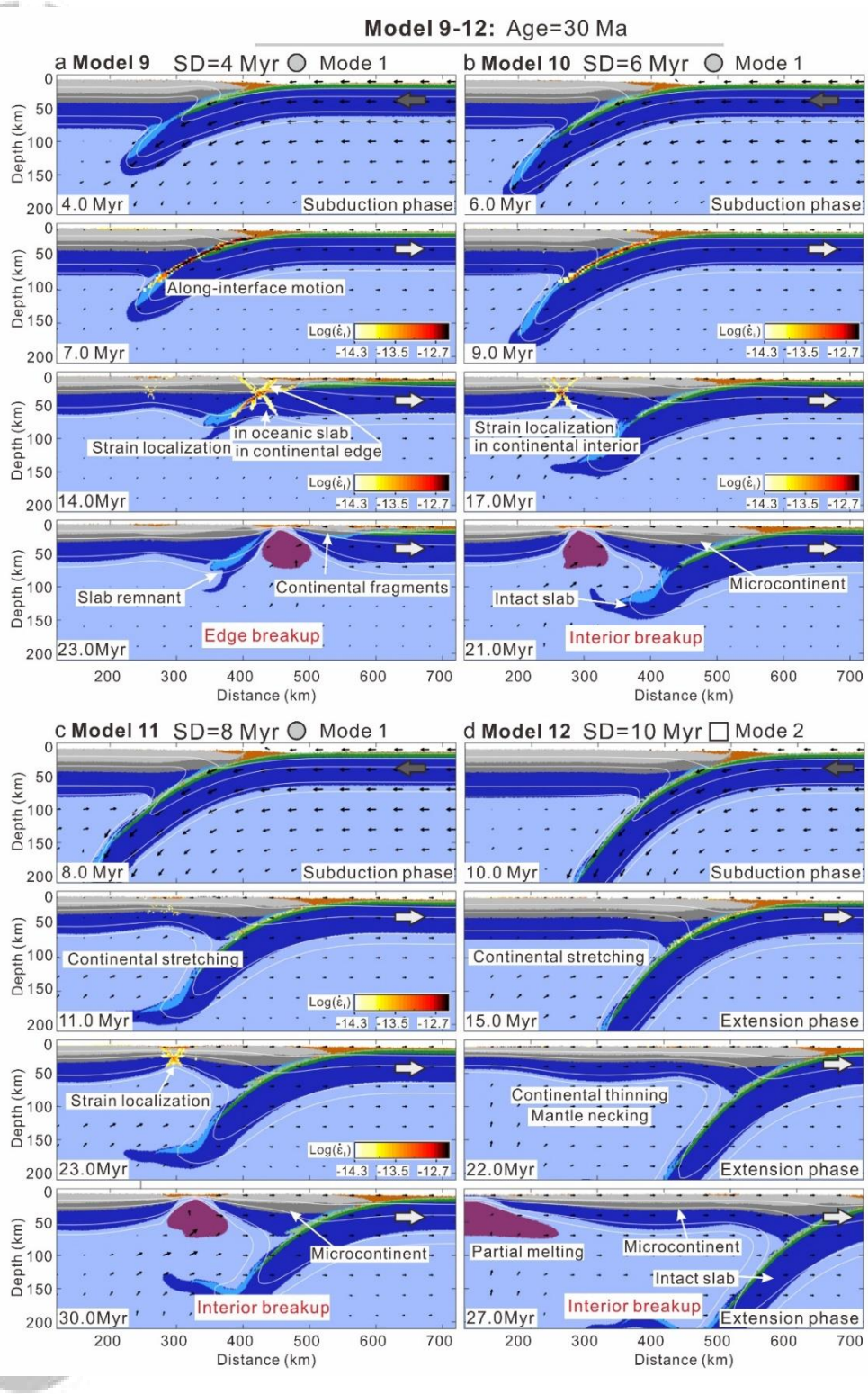

Figure 6. Evolution of group 3. Oceanic lithospheric age of 30 Ma and subduction duration (SD) of 4 (model 9), 6 (model 10), 8 (model 11), 10 Myr (model 12). The inset colors show the second invariant of the strain rate (as in the colorbar). White lines are isotherms in °C. Black arrows denote the calculated velocity field. Colors of rock types are as on Figure 2.

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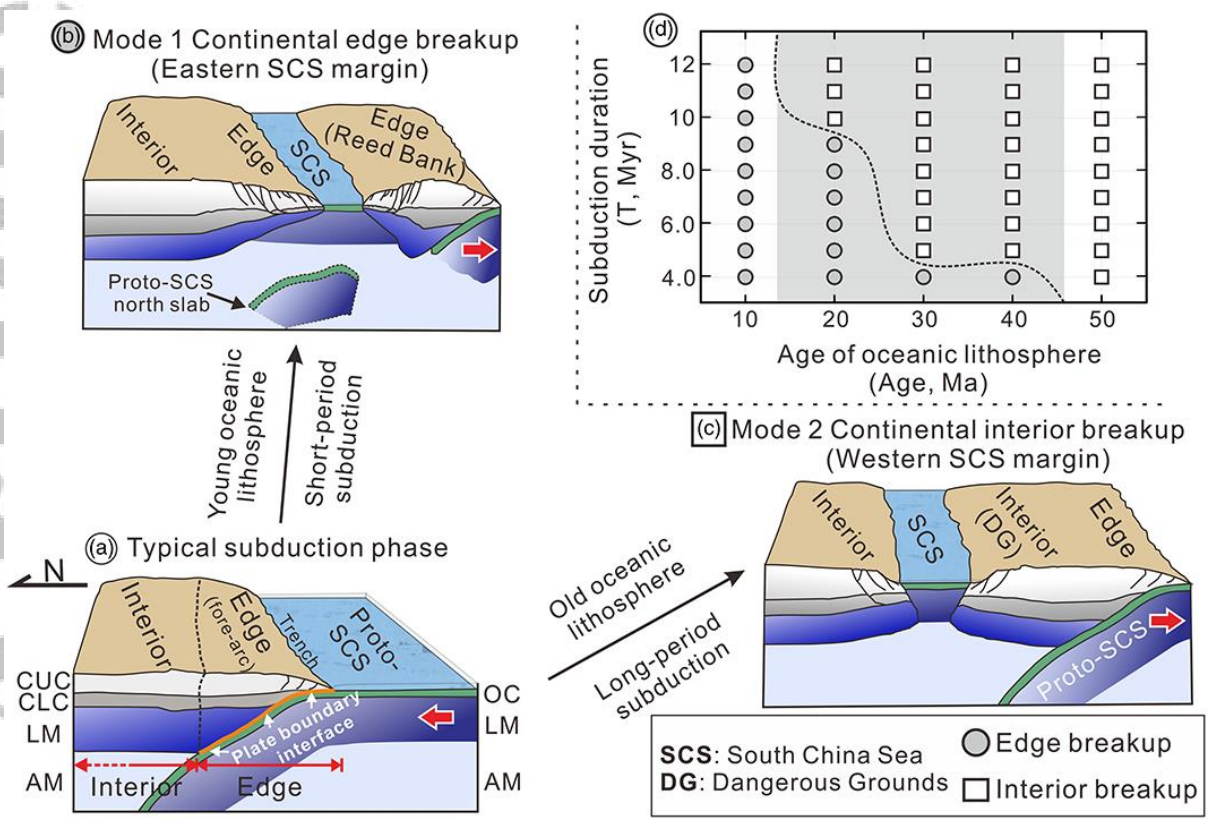

Figure 7. Conceptual models of the contrasting continental breakup modes (a-c) and area diagram summarizing the models in terms of breakup types in the function of subduction duration and age of the oceanic lithosphere (d). (a) Typical subduction zone, in which the upper plate is divided into continental edge and interior according to whether they are located directly above the slab-lithosphere interface. (b, c) Illustration of the breakup of the continental interior and edge, respectively. (d) The dashed line marks the threshold values indicating the transformation from continental edge to interior breakup. The shadow area denotes the conditions under which both breakup modes at a given age can be expected by varying the subduction duration. CUC: Continental upper crust; CLC: Continental lower crust; LM: Lithospheric mantle; AM: Asthenospheric mantle; OC: Oceanic crust.

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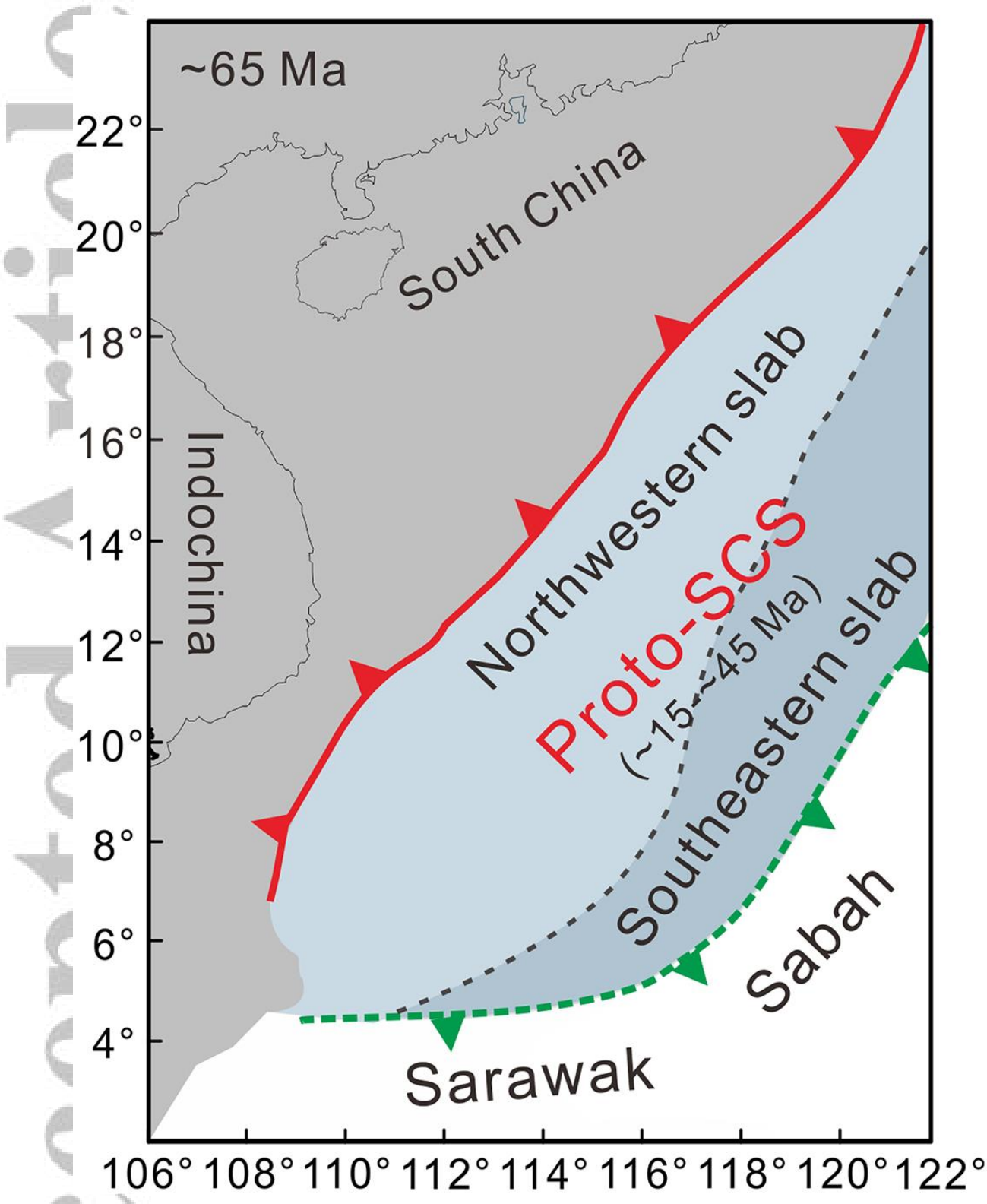

Figure 8. Sketch map showing the proportion of the Proto-SCS subducting towards the northwest and southeast. In the northwest dipping subduction (red line), a larger portion of the PSCS has been subducted beneath the western South China margin compared to that beneath

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the eastern South China margin. The dashed line divides the Proto-SCS into northwestern and southeastern slabs which have been subducted beneath the South China block (along the red line) and Sarawak/Sabah (along the green line), respectively.

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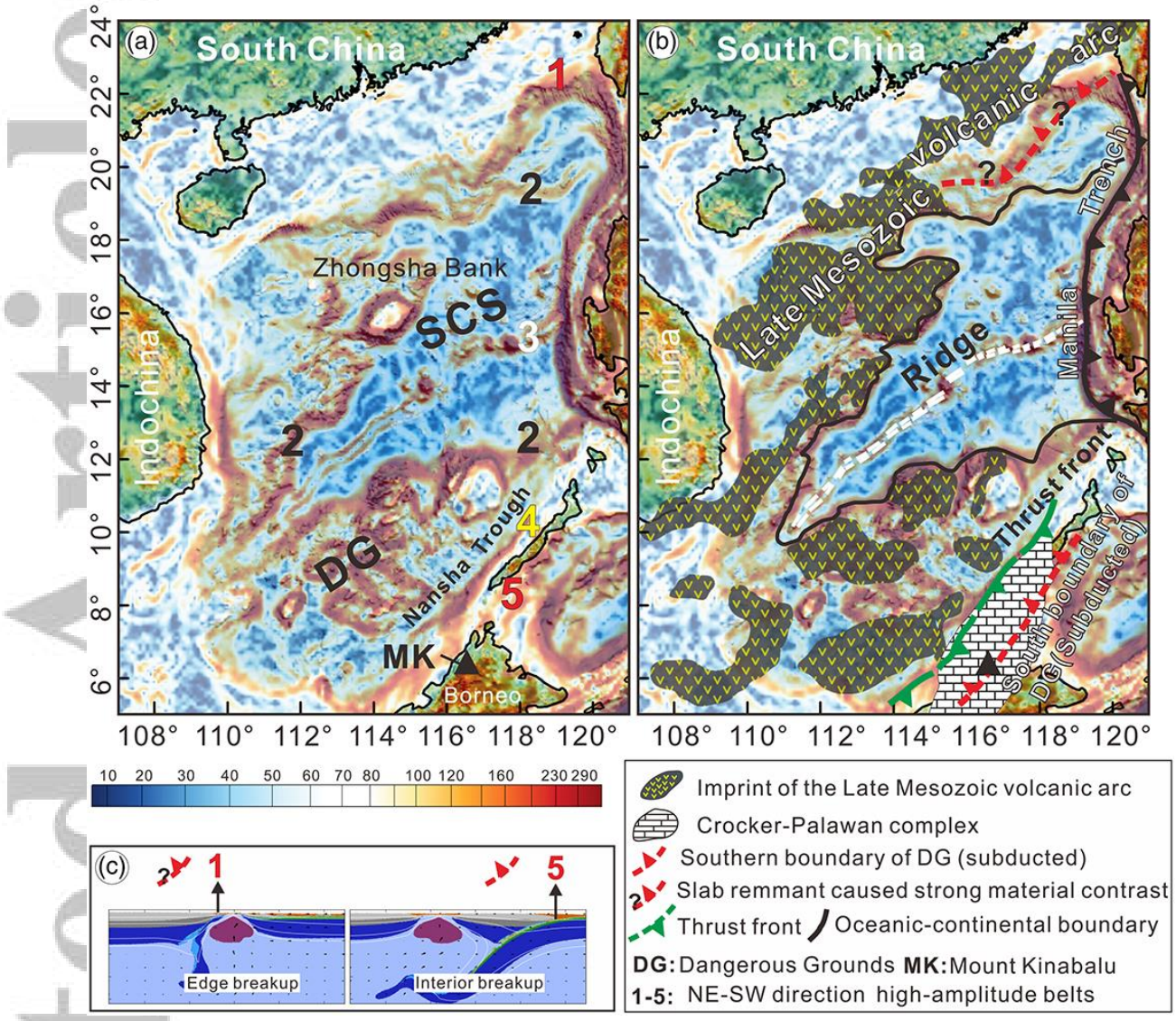

Figure 9. Map of the horizontal gradient of the Bouguer anomaly in the SCS region without (a) and with (b) interpretation, (c) the numerical models used for comparison. High-amplitude anomaly belts paralleling to the South China continental margin are numbered. Region of Crocker-Palawan complex is modified after *Hutchison, 2004*. The free-air gravity anomaly data are obtained from the satellite-derived field produced by *Sandwell et al., (2014)*. Terrain corrections for Bouguer gravity anomalies are derived from 1'x1' resolution ETOPO1 model (*Amante and Eakins, 2009*). Nansha Trough (NW Borneo trough), Nansha (Dangerous Grounds).

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